Lecture 2

Adsorption and the thermodynamics of surfaces Adsorption at gas liquid interface

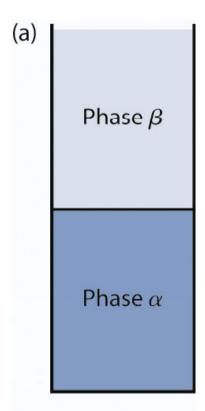
 Adsorption - a tendency of one component to have a higher or lower concentration at the interface in comparison to adjacent phase.

Models of the interface

 A system containing an interface can be divided into 3 regions

for an extensive property B:

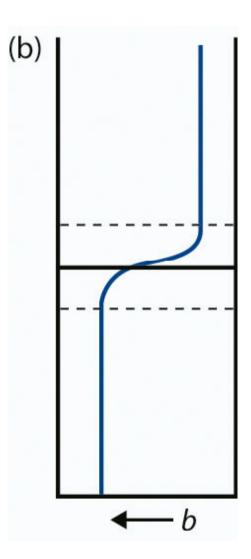
$$B = B^{\alpha} + B^{\beta} + B^{\sigma}$$



Models of the interface

1. Surface phase approach (Guggenheim)

- interfacial region has finite thickness;
- properties in the bulk phases are uniform



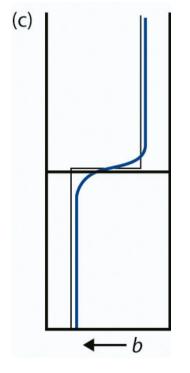
Models of the interface

2. Surface excess properties approach (Gibbs)

- interface is infinitely thin;
- properties of bulk phases assumed to extend till the interface;
- difference between the real property B and the model value is called excess

extensive property:
$$B_{\rm excess}=B^{\,\sigma}=B_{\rm real}-B_{\rm model}$$
 $V=V^{\,\alpha}+V^{\,\beta}\,, \quad V^{\,\sigma}=0$

intensive property: b = B/V

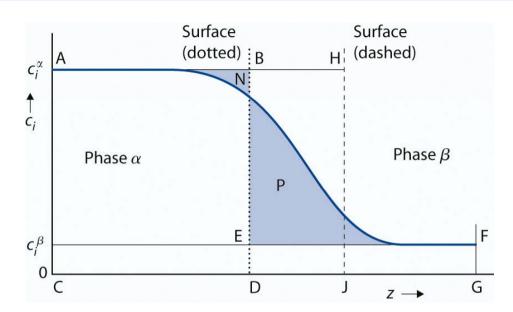


Example: concentration
$$n_{i, \text{model}} = c_i^{\ \alpha} V^{\ \alpha} + c_i^{\ \beta} V^{\ \beta}$$
 excess $n_i^{\ \sigma} = n_i - \left(c_i^{\ \alpha} V^{\ \alpha} + c_i^{\ \beta} V^{\ \beta}\right)$

surface excess:

$$n_i^{\ \sigma} = n_i - \left(c_i^{\ \alpha} V^{\alpha} + c_i^{\ \beta} V^{\beta}\right)$$

$$\Gamma_i = n_i^{\ \sigma} \ / \ A$$
 adsorption



the value of adsorption depends of the interface position!
 so, we need to agree on a convention....

Gibbs convention:

adsorption of a major component is zero:

$$\left|\Gamma_{A}=n_{A}^{\sigma}/A=0\right|$$

relative adsorption

$$n_i^{\sigma} = n_i - \left(c_i^{\alpha} V^{\alpha} + c_i^{\beta} V^{\beta}\right) = n_i - c_i^{\alpha} V + \left(c_i^{\alpha} - c_i^{\beta}\right) V^{\beta}$$

$$n_A^{\sigma} = n_A - c_A^{\alpha} V + \left(c_A^{\alpha} - c_A^{\beta}\right) V^{\beta}$$

in the combination below V^b will be eliminated:

$$n_{i}^{\sigma} - n_{A}^{\sigma} \left(\frac{c_{i}^{\alpha} - c_{i}^{\beta}}{c_{A}^{\alpha} - c_{A}^{\beta}} \right) = n_{i} - c_{i}^{\alpha} V - \left(n_{A} - c_{A}^{\alpha} V \right) \left(\frac{c_{i}^{\alpha} - c_{i}^{\beta}}{c_{A}^{\alpha} - c_{A}^{\beta}} \right)$$



doesn't depend on interface position!

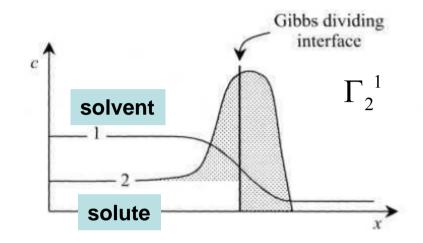
normalizing this by area will make relative adsorption

$$\Gamma_{i}^{A} = \Gamma_{i} - \Gamma_{A} \frac{c_{i}^{\alpha} - c_{i}^{\beta}}{c_{A}^{\alpha} - c_{A}^{\beta}}$$

$$\Gamma_{i}^{A} = \Gamma_{i} - \Gamma_{A} \frac{c_{i}^{\alpha} - c_{i}^{\beta}}{c_{A}^{\alpha} - c_{A}^{\beta}}$$

Gibbs convention:

$$\Gamma_A = 0, \ \Gamma_i^A = \Gamma_i(Gibbs)$$



 The relative adsorption of *i* with respect to A is equal to the adsorption of *i* using Gibbs convention

 <u>adsorption isotherm</u>: dependence of the adsorption vs. bulk concentration at constant temperature. Adsorption isotherm is described by an <u>isotherm equation</u>

- Consider a system of two phases separated by an interface
- Let's start with the internal energy as it only contains extensive quantities as variables:

Generally:
$$dU = \delta q + \delta w = TdS - PdV + \sum_{i} \mu_{i} dn_{i}$$

$$dU = dU^{\alpha} + dU^{\beta} + dU^{\sigma}$$

$$dU = TdS^{\alpha} - P^{\alpha} dV^{\alpha} + \sum_{i} \mu_{i}^{\alpha} dN_{i}^{\alpha}$$

$$TdS^{\beta} - P^{\beta} dV^{\beta} + \sum_{i} \mu_{i}^{\beta} dN_{i}^{\beta}$$

$$TdS^{\sigma} + \gamma dA + \sum_{i} \mu_{i}^{\sigma} dN_{i}^{\sigma}$$

Sum over all components present in the system

As the interface is infinitely thin, it cannot perform volume work.

Rearranging:

$$dU = TdS - P^{\alpha}dV - \left(P^{\beta} - P^{\alpha}\right)dV^{\beta} + \sum_{i} \mu_{i}^{\alpha}dN_{i}^{\alpha} + \sum_{i} \mu_{i}^{\beta}dN_{i}^{\beta} + \sum_{i} \mu_{i}^{\sigma}dN_{i}^{\sigma} + \gamma dA$$

for Helmholtz energy:

$$dF = -SdT - P^{\alpha}dV - \left(P^{\beta} - P^{\alpha}\right)dV^{\beta} + \sum_{i} \mu_{i}^{\alpha}dN_{i}^{\alpha} + \sum_{i} \mu_{i}^{\beta}dN_{i}^{\beta} + \sum_{i} \mu_{i}^{\sigma}dN_{i}^{\sigma} + \gamma dA$$
 =0 at T, V = constant.

• in a closed system: $dN_i^{\sigma} = -(dN_i^{\alpha} + dN_i^{\beta})$

$$\frac{dF}{dN_i^{\alpha}} = \mu_i^{\alpha} - \mu_i^{\sigma} = 0; \frac{dF}{dN_i^{\beta}} = \mu_i^{\beta} - \mu_i^{\sigma} = 0$$

$$dF = \left(P^{\beta} - P^{\alpha}\right) dV^{\beta} + \gamma dA; \frac{\partial F}{\partial A}\Big|_{T, V, V^{\beta}, N_i} = \gamma$$

Young-Laplace equation:

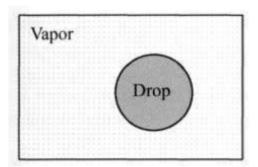
$$dF = -\left(P^{\beta} - P^{\alpha}\right)dV^{\beta} + \gamma dA;$$

$$\frac{dF}{dA} = \frac{\partial F}{\partial A} + \frac{\partial F}{\partial V^{\beta}} \frac{\partial V^{\beta}}{\partial A} = \gamma - \left(P^{\beta} - P^{\alpha}\right) \left(\frac{1}{R_{1}} + \frac{1}{R_{2}}\right)^{-1} = 0$$

For a case of a spherical droplet:

$$P^{\beta} - P^{\alpha} = 2\gamma/r = 2\gamma'/r'$$

Surface tension depends on the location of Gibbs plane in a case of curved interface! No such problems on a planar interface...



Similarly for Gibbs free energy:

$$dG = -SdT + V^{\alpha}dP^{\alpha} + V^{\beta}dP^{\beta} + \gamma dA + \sum_{i} \mu_{i}dn_{i}$$

in a case of planar interface:

$$dG = -SdT + VdP + \gamma dA + \sum_{i} \mu_{i} dn_{i}$$

$$\left. \frac{\partial G}{\partial A} \right|_{T,P,N_i} = \gamma$$

 at constant temperature we can integrate the surface term and normalize it by the area

$$\frac{G^{\sigma}}{A} = \gamma + \sum_{i} \mu_{i}^{\sigma} d\Gamma_{i}^{\sigma}$$

$$\frac{U^{\sigma}}{A} = \frac{H^{\sigma}}{A} = \frac{TS^{\sigma}}{A} + \gamma + \sum_{i} \mu_{i}^{\sigma} d\Gamma_{i}^{\sigma}$$

Surface excess of thermodynamic functions:

$$dU^{\sigma} = TdS^{\sigma} + \gamma dA + \sum_{i} \mu_{i} dN_{i}^{\sigma}$$
$$dF^{\sigma} = -S^{\sigma} dT + \gamma dA + \sum_{i} \mu_{i} dN_{i}^{\sigma}$$

• As U is a linear homogeneous function:

$$U^{\sigma} = TS^{\sigma} + \gamma A + \sum_{i} \mu_{i} N_{i}^{\sigma}$$

$$F^{\sigma} = U^{\sigma} - TS^{\sigma} = \gamma A + \sum_{i} \mu_{i} N_{i}^{\sigma}$$

- in a case of planar interface:
- entropy-surface tension relation:

Maxwell relations for exact differential:

$$f = g dx + h dy \implies \frac{\partial g}{\partial y}\bigg|_{x} = \frac{\partial h}{\partial x}\bigg|_{y}$$
 heat per unit area
$$-\frac{\partial S^{\sigma}}{\partial A}\bigg|_{T,N_{i}} = \frac{\partial \gamma}{\partial T}\bigg|_{A,N_{i}} \implies s^{\sigma} = -\frac{\partial \gamma}{\partial T}\bigg|_{A,N_{i}} \quad u^{\sigma} = \gamma - T\frac{\partial \gamma}{\partial T}\bigg|_{A,P}; \ q = -T\frac{\partial \gamma}{\partial T}$$

Gibbs adsorption isotherm

Let's consider internal energy:

$$dU^{\sigma} = TdS^{\sigma} + \gamma dA + \sum_{i} \mu_{i} dN_{i}^{\sigma}$$

$$U^{\sigma} = TS^{\sigma} + \gamma A + \sum_{i} \mu_{i} N_{i}^{\sigma}$$

Gibbs-Duhem equation:

$$S^{\sigma}dT + Ad\gamma + \sum_{i} N_{i}^{\sigma}d\mu_{i} = 0$$

at constant temperature:

$$d\gamma = -\sum_{i} \Gamma_{i}^{\sigma} d\mu_{i}$$

Gibbs Adsorption Isotherm

for 2-component mixture:

$$\Gamma_2^{(1)} = -\frac{1}{RT} \frac{d\gamma}{d(\ln a)} = -\frac{a}{RT} \frac{d\gamma}{da} \qquad \text{if } \Gamma > 0 \quad d\gamma/da < 0 \\ \text{if } \Gamma < 0 \quad d\gamma/da > 0$$

Marangoni effect

Tears of wine



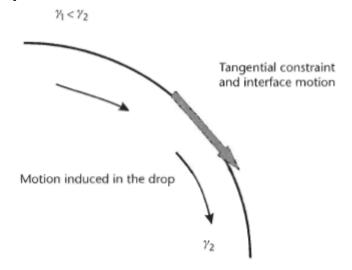
Marangoni effect

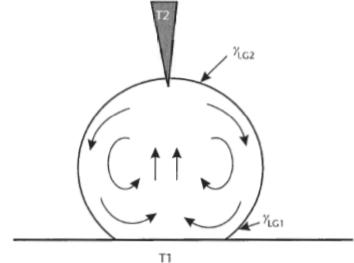
surface tension depends on temperature as

$$\gamma = \gamma_0 (1 - \beta (T - T_0))$$

for water/air interface: γ_0 =72mN/m and β =0.1 mN/(m K)

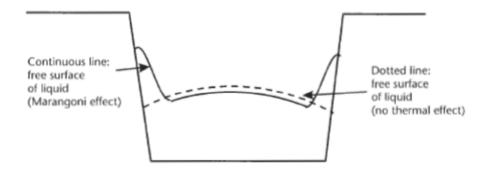
 surface tension distribution induces tangentional force distribution on the interface and convective motion inside the droplet



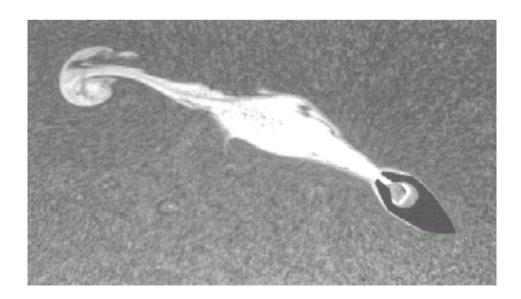


Marangoni effect

Marangoni effect due to temperature in a microwell



Marangoni effect due to surfactant concentration



Measurement of adsorption

Methods for measuring adsorption can be grouped into following categories

adsorption from concentration change

$$n_i^{\sigma} = n_i^{0} - n_i = (c_i^{0} - c_i)V$$

adsorption from surface analysis (sampling at the surface)

$$n_{i,real} = c_i^L V^L + c_i^S V^S + c_i^G V^G$$

$$n_i^{\sigma} = n_{i,real} - n_{i,model} = \left(c_i^S - c_i^L\right) V^S$$

adsorption from surface tension change

Gibbs equation:
$$\frac{-d\gamma}{RT} = \sum_{i} \left(\Gamma_{i} d \ln(a_{i}) \right)$$

for 2 component system:
$$\Gamma_B = \frac{a_B}{RT} \frac{d\gamma}{d(a_B)}$$

Measurement of adsorption

adsorption from concentration change

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adsorption from surface analysis (sampling at the surface)

$$n_{i,real} = c_i^L V^L + c_i^S V^S + c_i^G V^G$$

$$n_i^{\sigma} = n_{i,real} - n_{i,model} = \left(c_i^S - c_i^L\right) V^S$$

adsorption from surface tension change

Gibbs equation:
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for 2 component system:
$$\Gamma_B = \frac{a_B}{RT} \frac{d\gamma}{d(a_B)}$$

Adsorption at Gas-Liquid interface

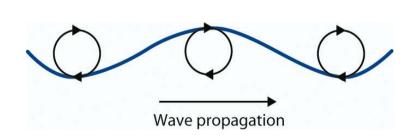
- Measurements of equilibrium adsorption
 - surface tension measurements (Wilhelmy plate)
 - surface analysis
 - radio-labelled solutes
 - neutron reflectometry (deuterated solutes)
 - X-ray reflectometry
 - formation and collection of foam

Adsorption at Gas-Liquid interface

Observation of adsorption kinetics

adsorption at freshly formed interfaces

- surface waves:
 - transverse capillary waves (ripples generated by an oscillating hydrophobic knife edge, f=100-300 Hz). Energy dissipation caused by compressionexpansion cycle.
 - longitudinal waves (horizontal movement of barrier, <0.1 Hz)



Wilhelmy plate

Overflow

Sintered glass

Pump

Adsorption of non-electrolyte solutes

Negative adsorption

observed in dilute aqueous solutions of e.g. glycine and sucrose

Kinetics of adsorbtion

in case of no stirring and no energy barrier:

$$\frac{d\Gamma}{dt} = \left(\frac{D}{\pi}\right)^{1/2} ct^{-1/2}, \quad \Gamma = 2ct \left(\frac{D}{\pi}\right)^{1/2}$$

Adsorption of ionized solute

 in case ionized solute the Gibbs equation must include contribution of all ions:

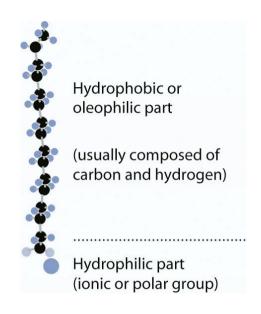
from electro neutrality:

$$\frac{-d\gamma}{RT} = \Gamma_{M^+} d \ln c_{M^+} + \Gamma_{X^-} d \ln c_{X^-}$$

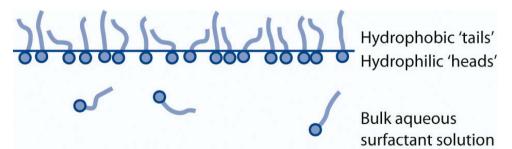
$$c_{M^+} = c_{X^-} = c \text{ and } \Gamma_{M^+} = \Gamma_{X^-} = \Gamma$$

$$\Gamma = \frac{1}{2RT} \frac{-d\gamma}{dc}$$

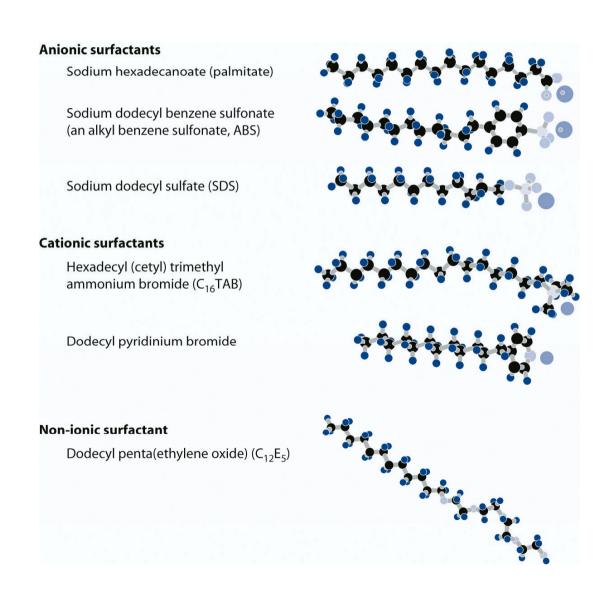
 surfactants (stands for: surface active agents) belong to a class of amphiphiles



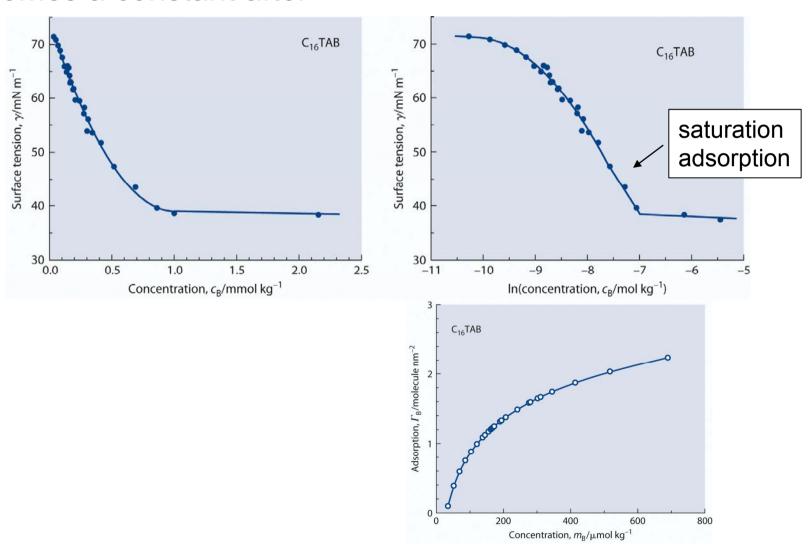
Gibbs monolayers



- Surfactants can be:
 - anionic,
 - cationic,
 - non-ionic



 Surface tension of surfactant usually falls to a lower limit and becomes a constant after



- Gibbs equation occasionaly reveals discrepancies with experimentally measured values.
 - surface hydrolysis: one of the ions is not adsorbed by the surface and replaced by H⁺.

$$-\frac{d\gamma}{RT} = \Gamma_{M^{+}} d \ln(c_{M^{+}}) + \Gamma_{X^{-}} d \ln(c_{X^{-}}) + \Gamma_{H^{+}} d \ln(c_{H^{+}})$$

$$= 0$$

$$-\frac{d\gamma}{RT} = \Gamma_{X^{-}} d \ln(c_{X^{-}})$$

 indifferent electrolyte: same situation can be created artificially to avoid ambiguities that might arise due to surface hydrolysis by adding large amounts of M+Y-:

$$-\frac{d\gamma}{RT} = \Gamma_{M^{+}} d \ln(c_{M^{+}}) + \Gamma_{X^{-}} d \ln(c_{X^{-}}) + \Gamma_{Y^{-}} d \ln(c_{Y^{-}}) \qquad \Longrightarrow \qquad -\frac{d\gamma}{RT} = \Gamma_{X^{-}} d \ln(c_{X^{-}})$$

$$= \text{const}$$

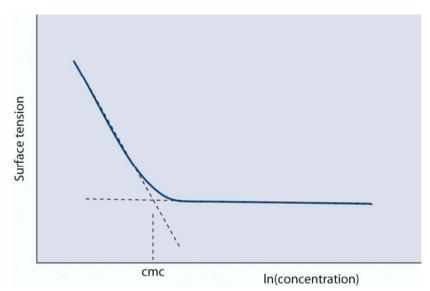
Micelles

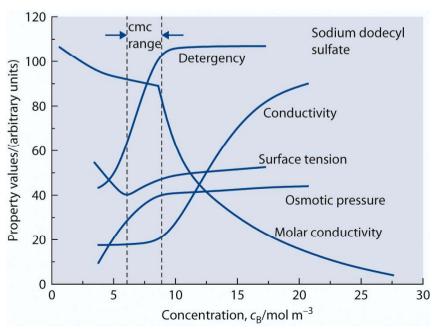
 above certain critical concentration the surface tension becomes independent of concentration: critical micelle concentration (cmc).

$$-\frac{d\gamma}{RT} = \Gamma_{X_S} d \ln(c_{X_S}) + \Gamma_{X_M} d \ln(c_{X_M})$$

not all amphiphiles form micelles.

- at cmc other properties show distinct changes as well (e.g. osmotic pressure indicate that the number of "solute particles" stays the same above cmc).
- formation of micelles means decrease in G, primarily due to large increase in entropy: hydrophobic interaction.





Probelms

• Problems 1, 2, 3